

Thrust and Efficiency of Propulsion by Oscillating Foils

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1 Introduction

Oscillating foils are being considered for propulsion of submersible vehicles and micro air vehicles (MAVs), at low Reynolds numbers where conventional propulsion techniques become inefficient. Linearized theories and potential flow methods provide good estimates of thrust and propulsive efficiencies at higher Reynolds numbers, but cannot account for the leading-edge flow separation and vortex formation that can occur in viscous flows. A 2D unsteady Navier Stokes solver was used to explore the effect of flow separation on the thrust and efficiency of an airfoil undergoing various combinations of sinusoidal pitching and plunging motions.

The Navier Stokes solver has been validated against a number of flow visualizations for pure plunging motion, and quantitative experimental and computational studies of pure pitching motion, with good results [1]. Here combined pitching and plunging motions are computed, with thrust and efficiency results compared to small-amplitude potential theory (Garrick [2]), large-amplitude unsteady panel method (UPM) potential flow simulations [1, 3], and experimental results of Anderson et al [4].

2 Numerical Method

The unsteady flow field around the oscillating airfoil is calculated using a 2nd-order accurate unsteady 2D compressible Navier Stokes solver. Motion of the airfoil is introduced by a combination of rigid-body motion and deformation of the grid, resulting in time-dependent grid metrics and Jacobian. Details of the method may be found in Young and Lai [1, 3], and Tuncer and Platzer [5]. Versions of the code have been parallelized with OpenMP for use on an Alpha cluster, and Parallel Virtual Machine (PVM) for use on a cluster of Pentium-based workstations running Linux.

The motion of the airfoil is a combination of plunging, $y = a \sin(\omega t)$ and pitching, $\theta = \theta_0 \sin(\omega t + \phi)$. Parameters of interest are the reduced frequency

$k = \omega c/2U_\infty$, non-dimensionalized plunge amplitude $h = a/c$, pitch amplitude θ_0 , pivot point of the pitching motion, and the phase angle ϕ (c is the airfoil chord and U_∞ is the free-stream velocity). The Strouhal number is calculated based on the frequency of motion and the total (peak to peak) excursion A_{TE} of the trailing edge, $St = \omega A_{TE}/2\pi U_\infty = kA_{TE}/\pi c$. Outputs include the time-averaged thrust coefficient $\overline{C_T} = -1/T \int_t^{t+T} C_D(t) dt$, input power coefficient $\overline{C_{P_{in}}} = -1/T \int_t^{t+T} [C_L(t)\dot{y}(t)/c + C_M(t)\dot{\theta}(t)] dt$, and propulsive efficiency, $\eta = \overline{C_T} U_\infty / \overline{C_{P_{in}}}$. C_D , C_L and C_M are the drag, lift and moment coefficients calculated by integrating viscous and pressure forces around the airfoil surface.

Optimization of time-averaged thrust and propulsive efficiency is performed using a gradient-descent based method. The objective function is a linear combination of the thrust coefficient and the efficiency, and the optimization variables are h , θ_0 , k , and ϕ . Evaluation of the gradient vector involves unsteady flow computation over a number of motion cycles until periodic behavior is established. Each of the gradient vector components is calculated in parallel. Calculations were typically performed on a 541×61 grid, with 189 points on the airfoil, a first grid point at 8.75×10^{-6} chords from the airfoil surface, and a boundary 20 chords from the airfoil in all directions.

3 Comparison of Numerical and Experimental Results

Figs. 1, 2, and 3 show the time-averaged thrust coefficient and the propulsive efficiency vs Strouhal number for three different motions. The parameter $\alpha_0 = -\tan^{-1}(2kh) + \theta_0$ is kept constant at 15° for all cases. Reynolds number $Re = 40,000$, and the airfoil is pitching about a point $1/3$ chords aft of the leading edge. These conditions were chosen to match those of the experiment performed by Anderson et al [4]. Navier Stokes calculations were performed using fully laminar flow, fully turbulent flow with the Baldwin-Lomax turbulence model, and a transition model using the Spalart-Allmaras turbulence model. Negligible differences in thrust and efficiency were found between the Baldwin-Lomax and Spalart-Allmaras models. Fig. 1, for which leading edge separation is minimal across the range of Strouhal number, shows relatively good agreement with the measurements of Anderson et al [4], and little difference between laminar and turbulent flows. Figs. 2 and 3, for which large leading edge separation is apparent at some Strouhal numbers, show lower efficiencies than measured by Anderson et al, and indicate a high degree of sensitivity to the precise details of leading edge vortex formation, dependent on whether the flow is assumed laminar or turbulent. In all three cases the level of agreement with experiment is good for the thrust coefficient, and moderate for the propulsive efficiency although the trends are well predicted.

It should be noted that small discrepancies in time-varying quantities may become relatively larger errors in the time-averaged quantities.

4 Optimization of Thrust and Efficiency

Figs. 1, 2 and 3 taken together show that the propulsive efficiency of an oscillating airfoil is sensitively dependent on the phase angle between pitching and plunging motion. This is in agreement with the findings reported by Isogai et al [6] for $Re = 10^5$ and Tuncer and Kaya [7], suggesting that high efficiencies and thrusts may be simultaneously achieved by optimizing interaction of the leading edge vortex with the foil motion.

The optimization procedure was checked by allowing the frequency and pitch amplitude to vary while holding other parameters fixed, and optimizing against efficiency. The results are shown in Fig. 2. In both laminar and turbulent cases the optimized solution converges to a point close to the peak efficiency found by a parametric search. In Table 1 the plunge amplitude is held constant at $h = 0.75$, $\alpha_0 = 15^\circ$, and the oscillation frequency and pitch amplitude are optimized. The procedure is repeated both with phase angle held constant, and optimized. Laminar and turbulent flows are considered. The table shows that the optimum phase angle between pitch and plunge motion is somewhere between $\phi = 75^\circ$ to 85° , again in good agreement with Anderson et al [4] and Isogai et al [6].

The effect of leading edge separation on thrust and efficiency may be seen in Figs. 4 and 5. Here numerical particle traces are shown for four points in the motion cycle for turbulent flow at two different Strouhal numbers. Fig. 4 corresponds to the highest efficiency point of Fig. 3 ($St = 0.184$, $\overline{C_T} = 0.24$, $\eta = 0.72$), and Fig. 5 is a high thrust, low efficiency case ($St = 0.70$, $\overline{C_T} = 1.22$, $\eta = 0.44$). Optimization of the flapping motion for highest efficiency appears to minimize leading edge separation. Optimization for high thrust promotes leading edge separation on the airfoil surface which is tilted upstream, so the low pressure of the separated vortex results in thrust rather than drag. The low thrust, low efficiency cases in Figs. 2 and 3 correspond to leading edge separation on the airfoil surface that is tilted downstream.

5 Conclusion

Simulations of the flow over an oscillating airfoil show good agreement with the experimental results of Anderson et al [4], for a range of different flapping parameters. The propulsive efficiency peaks at a Strouhal number between 0.15-0.3, as found by Anderson et al and other researchers. Gradient-descent based optimization of the phase angle between the pitching and plunging components of the motion shows a phase angle between $\phi = 75^\circ$ and 85° for best propulsive efficiency, in agreement with Anderson et al [4] and Isogai et

al [6]. Highest efficiency propulsion occurs when leading edge separation is avoided or minimized, however thrust is low in this case. High thrust corresponds to long periods of high effective angle of attack, and large separated leading edge vortices.

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Case	St	k	θ_0	ϕ	η
Turbulent, ϕ fixed	0.244	0.97	21.1°	75.0°	0.722
Turbulent, ϕ fixed	0.198	0.80	16.1°	90.0°	0.705
Turbulent, ϕ optimized	0.252	1.00	21.9°	78.6°	0.726
Laminar, ϕ fixed	0.280	1.10	24.5°	75.0°	0.598
Laminar, ϕ fixed	0.287	1.12	25.0°	90.0°	0.636
Laminar, ϕ optimized	0.258	1.02	22.5°	83.3°	0.643

Table 1. Optimization results, $Re = 40,000$, $\alpha_0 = 15^\circ$, pivot at $1/3$ chord, $h = 0.75$.

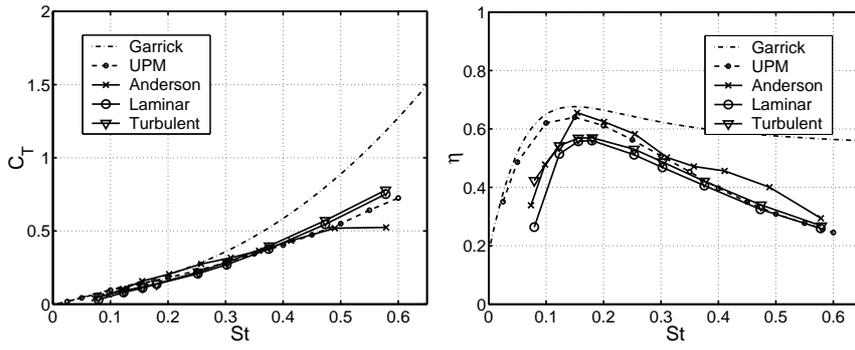


Fig. 1. Mean thrust coefficient and propulsive efficiency, $Re = 40,000$, $\alpha_0 = 15^\circ$, pivot at $1/3$ chord, $h = 0.25$, $\phi = 90^\circ$.

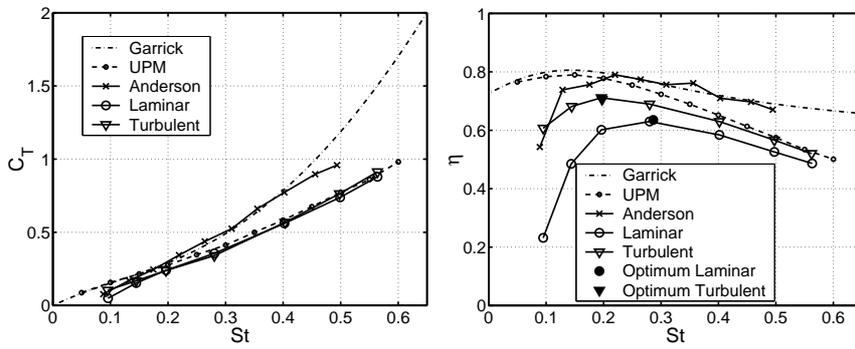


Fig. 2. Mean thrust coefficient and propulsive efficiency, $Re = 40,000$, $\alpha_0 = 15^\circ$, pivot at $1/3$ chord, $h = 0.75$, $\phi = 90^\circ$.

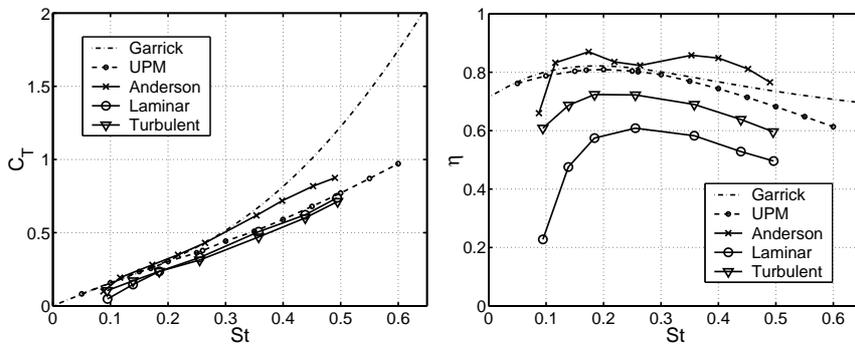


Fig. 3. Mean thrust coefficient and propulsive efficiency, $Re = 40,000$, $\alpha_0 = 15^\circ$, pivot at $1/3$ chord, $h = 0.75$, $\phi = 75^\circ$.

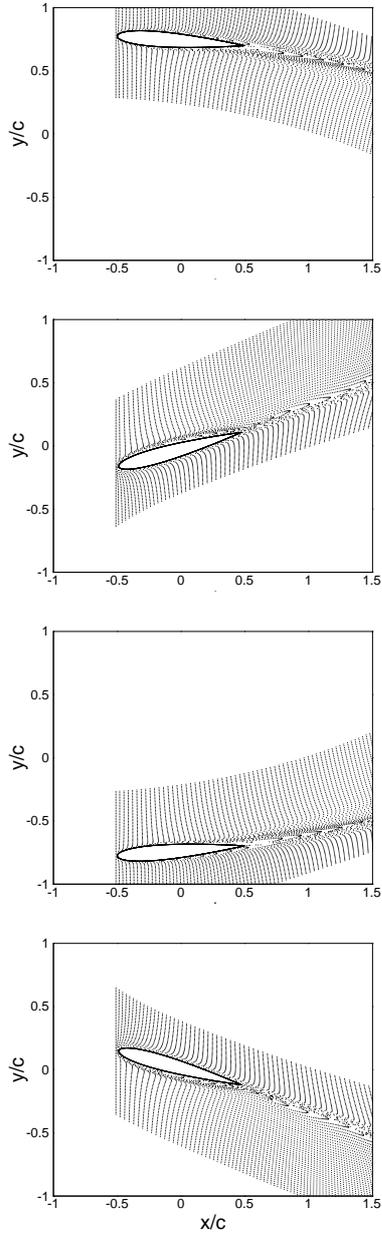


Fig. 4. Particle traces, $\alpha_0 = 15^\circ$, $h = 0.75$, $\phi = 75^\circ$, $St = 0.184$

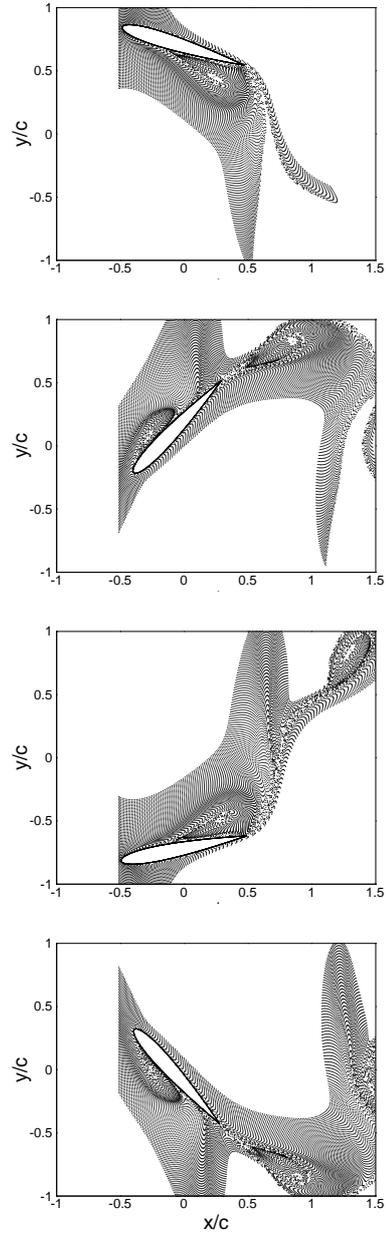


Fig. 5. Particle traces, $\alpha_0 = 15^\circ$, $h = 0.75$, $\phi = 75^\circ$, $St = 0.70$