

Electromagnetic interactions

$$\vec{F} = q \left(\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right)$$

$$H = T + V$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\dot{\vec{\Pi}} = \frac{d\vec{\Pi}}{dt}$$

$\vec{\Pi}$: mechanical momentum

$$\vec{\Pi} = m\vec{v} \quad (\text{classically})$$

$$\left. \begin{array}{l} \vec{A} \rightarrow \vec{A} + \nabla \Lambda \\ \phi \rightarrow \phi - \frac{1}{c} \frac{\partial \Lambda}{\partial t} \end{array} \right\} \begin{array}{l} \vec{B} \rightarrow \vec{B} \\ \vec{E} \rightarrow \vec{E} \end{array}$$

$$\frac{d\vec{\Pi}}{dt} = q \left[-\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} + \frac{\vec{v}}{c} \times (\nabla \times \vec{A}) \right]$$

$$\begin{aligned} \vec{v} \times (\nabla \times \vec{A}) &= (\nabla (\vec{v} \cdot \vec{A}) - (\nabla \cdot \vec{v}) \vec{A}) \\ &= \nabla (\vec{v} \cdot \vec{A}) - (\nabla \cdot \vec{v}) \vec{A} \end{aligned}$$

$$\frac{d\vec{\Pi}}{dt} = q \left\{ -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \frac{1}{c} \left[\nabla (\vec{v} \cdot \vec{A}) - (\nabla \cdot \vec{v}) \vec{A} \right] \right\}$$

$$\begin{aligned} \frac{d\vec{\Pi}}{dt} &= -\nabla \left\{ q \left(\phi - \frac{1}{c} \vec{v} \cdot \vec{A} \right) \right\} \\ &\quad - \frac{1}{c} q \left(\frac{\partial \vec{A}}{\partial t} + (\nabla \cdot \vec{v}) \vec{A} \right) \end{aligned}$$

$$\frac{d\vec{A}(\vec{r}(t), t)}{dt} = \frac{\partial \vec{A}}{\partial t} + \frac{\partial \vec{A}}{\partial x_i} \frac{\partial x_i}{\partial t}$$

$$\frac{d\vec{A}}{dt} = \frac{\partial \vec{A}}{\partial t} + (\vec{v} \cdot \nabla) \vec{A}$$

$$\frac{d\pi}{dt} = -\nabla \left[q \left(\phi - \frac{1}{c} \vec{v} \cdot \vec{A} \right) \right]$$

$$\frac{d}{dt} \left(\vec{\pi} + \frac{q}{c} \vec{A} \right) = -\nabla \left[q \left(\phi - \frac{1}{c} \vec{v} \cdot \vec{A} \right) \right]$$

$$\frac{d\vec{p}}{dt} = -\nabla V_{\text{eff}}$$

$$V_{\text{eff}} = q \left(\phi - \frac{1}{c} \vec{v} \cdot \vec{A} \right)$$

$$\vec{p} = \vec{\pi} + \frac{q}{c} \vec{A} \quad \text{: canonical momentum}$$

$$H = \frac{\vec{\pi}^2}{2m} + q\phi$$

$$H = \frac{\left(\vec{p} - \frac{q}{c} \vec{A} \right)^2}{2m} + q\phi$$

$$\begin{aligned}
\frac{dx}{dt} &= \frac{1}{i\hbar} [H, x] \\
&= \frac{1}{i\hbar} \frac{1}{2m} \left[\left(\vec{p} - \frac{q}{c} \vec{A} \right)^2, x \right] \\
&= \frac{1}{i\hbar} \frac{1}{2m} \left[\left(p_x - \frac{q}{c} A_x \right)^2, x \right] \\
&= \frac{1}{i\hbar} \frac{1}{2m} \left\{ \left(p_x - \frac{q}{c} A_x \right) \left[p_x - \frac{q}{c} A_x, x \right] \right. \\
&\quad \left. + \left[p_x - \frac{q}{c} A_x, x \right] \left(p_x - \frac{q}{c} A_x \right) \right\} \\
&= \frac{1}{i\hbar} \frac{1}{2m} \left(p_x - \frac{q}{c} A_x \right) 2i\hbar \\
\frac{dx}{dt} &= \frac{1}{m} \left(p_x - \frac{q}{c} A_x \right)
\end{aligned}$$

$$\boxed{m \frac{dx}{dt} = \Pi_x}$$

Particle in a magnetic field

$$\phi = 0 \quad \vec{A}', \vec{A} = \vec{A} + \vec{\nabla} \Lambda(\vec{r})$$

$$\underline{\phi = 0, \phi' = 0}$$

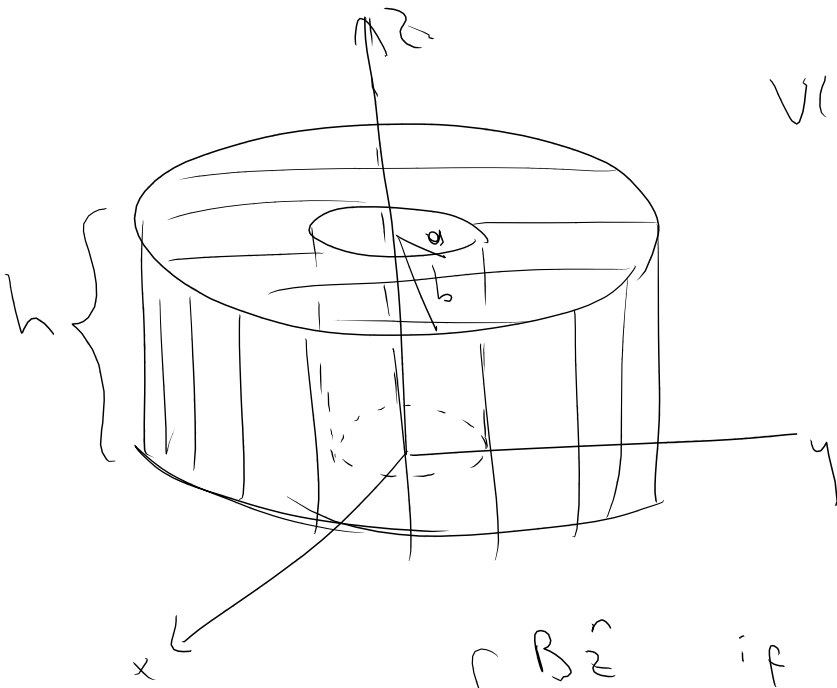
$$H = \frac{(\vec{p} - \frac{q}{c}\vec{A})^2}{2m} = \frac{1}{2m} \vec{p}^2 + \frac{1}{2m} \frac{q^2}{c^2} \vec{A}^2 - \frac{q}{2m c} (\vec{p} \cdot \vec{A} + \vec{A} \cdot \vec{p})$$

$$\langle \vec{r} | H | \alpha t_0; t \rangle = \langle \vec{r} | i\hbar \frac{\partial}{\partial t} | \alpha t_0; t \rangle$$

$$\langle \vec{r} | \frac{(\vec{p} - \frac{q}{c}\vec{A})^2}{2m} | \alpha t_0; t \rangle = i\hbar \frac{\partial}{\partial t} \psi(\vec{r}, t)$$

$$\frac{(-i\hbar \vec{\nabla} - \frac{q}{c}\vec{A}) \cdot (-i\hbar \vec{\nabla} - \frac{q}{c}\vec{A})}{2m} \langle \vec{r} | \vec{p} - \frac{q}{c}\vec{A} | \alpha t_0; t \rangle = i\hbar \frac{\partial}{\partial t} \psi(\vec{r}, t)$$

$$\left(-i\hbar \vec{\nabla} - \frac{q}{c}\vec{A} \right) \cdot \left(-i\hbar \vec{\nabla} - \frac{q}{c}\vec{A} \right) \psi(\vec{r}, t) = i\hbar \frac{\partial}{\partial t} \psi(\vec{r}, t)$$



$$V(\vec{r}) = \begin{cases} 0 & \text{if } 0 < z < h \\ & a < \rho < b \\ A & \text{otherwise} \end{cases}$$

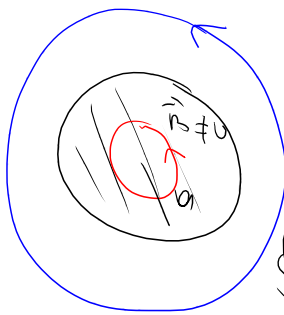
$$\vec{B}(\vec{r}) = \begin{cases} B\hat{z} & \text{if } \rho < a \\ 0 & \text{if } \rho > a \end{cases}$$

$$\vec{\nabla} \times \vec{A} = \vec{B} \quad \Rightarrow \quad \int d\vec{S} \cdot (\vec{\nabla} \times \vec{A}) = \int d\vec{S} \cdot \vec{B}$$

$$\vec{A} = \oint A(\rho)$$

$$\oint d\vec{l} \cdot \vec{A} = \Phi_B$$

$\vec{B} = 0$



$$\vec{B} = 0$$

$$d\vec{l} = dl \hat{\phi}$$

$$d\vec{l} \cdot \vec{A} = A(\rho) dl$$

$$\oint d\vec{l} \cdot \vec{A} = \oint A(\rho) dl = 2\pi \rho A(\rho)$$

$$\Phi_B = B \pi \rho^2$$

$$A(\rho) = \frac{B \pi \rho^2}{2\pi \rho} = \boxed{\frac{1}{2} B \rho = A(\rho)} \quad \rho < a$$

$$\Phi_B = B \pi a^2 \Rightarrow A(\rho) = \frac{B \pi a^2}{2\pi \rho} = \frac{B a^2}{2} \frac{1}{\rho} \quad \rho > a$$

$$\vec{A}(\rho) = \frac{B a^2}{2} \frac{1}{\rho} \hat{\phi} \quad \rho > a$$

$$\phi_0 = B a^2$$

$$-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}$$

$$= -i\hbar \left[\frac{\partial}{\partial \rho} \rho^n + \frac{\partial}{\partial \phi} + z^n \frac{\partial}{\partial z} \right] - \frac{q}{c} \frac{B a^2}{2\rho} \hat{\phi}$$

$$= -i\hbar \left[\rho^n \frac{\partial}{\partial \rho} + \frac{\partial}{\partial \phi} + \frac{1}{\rho} \left(\frac{\partial}{\partial \phi} - \frac{i}{\hbar} \frac{q}{c} \frac{\phi_0}{2\rho} \right) + z^n \frac{\partial}{\partial z} \right]$$

$$\left(-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A} \right)^2$$

$$= \frac{\hbar^2}{2} \left[\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \left(\frac{\partial}{\partial \phi} - \frac{i}{\hbar} \frac{q \phi_0}{2\rho} \right)^2 + \frac{\partial^2}{\partial z^2} \right]$$

$$\psi(\vec{r}) = R(\rho) \Phi(\phi) Z(z)$$

$$\psi(\vec{r})|_{z=0} = 0$$

$$\psi(\vec{r})|_{z=h} = 0$$

$$Z(0) = Z(h) = 0$$

$$\psi(\vec{r})|_{\rho=a} = 0$$

$$\psi(\vec{r})|_{\rho=b} = 0$$

$$R(a) = R(b) = 0$$

$$-\frac{\hbar^2}{2m} \left[\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \left(\frac{\partial}{\partial \phi} - \frac{i}{\hbar} \frac{q \Phi_0}{2\pi c} \right)^2 + \frac{\partial^2}{\partial z^2} \right] \Psi = E \Psi$$

$$E = -\frac{\hbar^2}{2m} \left[\frac{1}{R} \frac{1}{\rho} \frac{d}{d\rho} (\rho R') + \frac{1}{\rho^2} \frac{1}{\Phi} \left(\frac{d}{d\phi} - \frac{i}{\hbar} \frac{q \Phi_0}{2\pi c} \right)^2 \Phi + \frac{\partial^2}{\partial z^2} \right] \Psi$$

$\rho = kh$

$$\frac{\partial^2}{\partial z^2} = \text{const} = -k^2 \Rightarrow Z(z) = A \sin(kz)$$

$kh = n\pi \quad n = 1, 2, 3, \dots$

$$k_n = n \frac{\pi}{h}$$

$Z_n(z) = A_n \sin(k_n z)$

$$\left(E - \frac{\hbar^2 k_n^2}{2m} \right) = -\frac{\hbar^2}{2m} \left[\frac{1}{R} \frac{1}{\rho} \frac{d}{d\rho} (\rho R') \right] - \frac{\hbar^2}{2m} \frac{1}{\rho^2} \left\{ \frac{1}{\Phi} \left(\frac{d}{d\phi} - \frac{i}{\hbar} \frac{q \Phi_0}{2\pi c} \right)^2 \Phi \right\}$$

$$\frac{1}{\Phi} \left(\frac{d}{d\phi} - \frac{i}{\hbar} \frac{q \Phi_0}{2\pi c} \right)^2 \Phi = \text{const} = -\alpha^2$$

$$\Phi(\phi) = B e^{i m \phi} \quad m = 0, \pm 1, \pm 2, \pm 3, \dots$$

$$-\alpha^2 = \frac{1}{\cancel{Be^{i\alpha d}}} \left(im - \frac{i}{\hbar} \frac{q \phi_0}{2ac} \right)^2 \cancel{Be^{i\alpha d}}$$

$$\boxed{\alpha_m^2 = \left(m - \frac{q \phi_0}{2ac \hbar} \right)^2}$$

$$E - \frac{\hbar^2 k_n^2}{2m} = \frac{-\hbar^2}{2m} \left[\frac{1}{R} \frac{1}{\rho} \frac{d^2}{\rho} (\rho R') - \frac{1}{\rho^2} \alpha^2 \right]$$

$$\left[\frac{1}{R} \frac{1}{\rho} (\rho R'' + R') - \frac{1}{\rho^2} \alpha^2 \right] \frac{\hbar^2}{2m} + \left(E - \frac{\hbar^2 k_n^2}{2m} \right) = 0$$

$$\rho^2 R'' + \frac{2m}{\hbar^2} \left(E - \frac{\hbar^2 k_n^2}{2m} \right) \rho^2 R - \alpha_m^2 R = 0$$

\uparrow \uparrow \uparrow
 $\rho^2 R'$

$$H = \frac{1}{2m} \left(\vec{p} - \frac{q}{c} \vec{A} \right)^2$$

$$H \psi = E \psi$$

$$\vec{A} \rightarrow \vec{A}' = \vec{A} + \vec{\nabla} \Lambda$$

$$H \rightarrow H'$$

$$\psi_E \rightarrow \psi'_E$$

$$\psi'_E(\vec{r}) = U(\vec{r}) \psi_E(\vec{r})$$

$$\begin{aligned} & (-i\hbar \vec{\nabla} - \vec{A}') U \psi \\ &= U (-i\hbar \vec{\nabla} - \vec{A}) \psi \end{aligned}$$

$$(-i\hbar \vec{\nabla} - \vec{A}) \psi(\vec{r}) \rightarrow (-i\hbar \vec{\nabla} - \vec{A}') U \cdot \psi(\vec{r})$$

$$\Rightarrow U(\vec{r}) (-i\hbar \vec{\nabla} - \vec{A}) \psi(\vec{r})$$

$$H' \psi'_E = \frac{1}{2m} (-i\hbar \vec{\nabla} - \vec{A}')^2 \psi'_E(\vec{r})$$

$$= \frac{1}{2m} (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}') \cdot (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}') \psi'$$

$$= \frac{1}{2m} (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}') U (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}) \psi$$

$$= \frac{1}{2m} U (-i\hbar \vec{\nabla} - \vec{A})^2 \psi$$

$$= U H \psi$$

$$= U E \psi$$

$$H \psi'_E = E \psi'$$

$$\Rightarrow (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A}') U_n = U_n (-i\hbar \vec{\nabla} - \frac{q}{c} \vec{A})$$

$$\Rightarrow \boxed{\vec{A}' = \vec{A} + \vec{\nabla} \Lambda}$$

$$\vec{\nabla} (U_n^2) = U_n (\vec{\nabla}^2 U_n) + (\vec{\nabla} U_n)^2$$

$$\vec{\nabla} U_n = U_n \vec{\nabla} + (\vec{\nabla} U_n)$$

$$\begin{aligned} \cancel{-i\hbar U_n \vec{\nabla} - i\hbar (\vec{\nabla} U_n) - \frac{q}{c} \vec{A} U_n} &= \cancel{-i\hbar U_n \vec{\nabla} - \frac{q}{c} U_n \vec{A}} + \frac{q}{c} (\vec{\nabla} \Lambda) U_n \\ &= -i\hbar U_n \vec{\nabla} - \frac{q}{c} U_n \vec{A} \end{aligned}$$

$$\Rightarrow -i\hbar (\vec{\nabla} U_n) = \frac{q}{c} (\vec{\nabla} \Lambda) U_n$$

$$\boxed{U_n = e^{\frac{q}{c\hbar} \Lambda}}$$

$$\phi = 0$$

$$\vec{A} \rightarrow \vec{A} + \vec{\nabla} \Lambda$$

$$\psi \rightarrow e^{\frac{q}{c\hbar} \Lambda} \psi$$

$$\langle \vec{p} \rangle = \int d^3r \psi^*(\vec{r}) (-i\hbar \vec{\nabla}) \psi$$

$$\Rightarrow \int d^3r \psi^* e^{-\frac{q}{c\hbar} \Lambda} (-i\hbar \vec{\nabla}) e^{\frac{q}{c\hbar} \Lambda} \psi$$

$$= \int d^3r \psi^* e^{-\frac{q}{c\hbar} \Lambda} \left[\frac{q}{c} (\vec{\nabla} \Lambda) - i\hbar \vec{\nabla} \right] \psi$$

$$\langle \vec{p} \rangle = \langle \vec{\nabla} \Lambda \rangle \frac{q}{c} + \langle \vec{p} \rangle$$

$$\vec{\Pi} = \vec{p} - \frac{q}{c} \vec{A}$$

$$\begin{aligned} \langle \vec{\Pi} \rangle &= \langle \vec{p} \rangle - \frac{q}{c} \langle \vec{A} \rangle \\ &\Rightarrow \langle \vec{\nabla} \lambda \rangle \frac{q}{c} + \langle \vec{p} \rangle \\ &\quad - \frac{q}{c} (\langle \vec{A} \rangle + \langle \vec{\nabla} \lambda \rangle) \\ &= \langle \vec{p} - \frac{q}{c} \vec{A} \rangle \end{aligned}$$

$$\langle \vec{\Pi} \rangle \Rightarrow \langle \vec{\pi} \rangle$$

$$\vec{\Pi} = m \frac{d\vec{x}}{dt} \text{ in Heisenberg picture.}$$

$$g(\vec{r}) = \sqrt{\psi^* \psi} \rightarrow \sqrt{\psi^* \underbrace{U^\dagger U}_1 \psi} = \sqrt{\psi^* \psi}$$

$g(\vec{r})$ is Gauge invariant.

$\vec{j}(\vec{r}) \propto [\psi^* \nabla \psi]$ is not gauge invariant.

\Rightarrow is valid only if
 $H = \frac{p^2}{2m} + V(\vec{r})$

Now

$$H = \frac{1}{2m} \left(\vec{p} - \frac{q}{c} \vec{A} \right)^2 \quad \frac{\partial g}{\partial t} + \vec{\nabla} \cdot \vec{j} = 0$$

$$\langle \vec{r} | \hat{H} | \vec{r} \rangle = \int d^3r \psi^* \left(\frac{\vec{p}}{m} - \frac{q}{c} \vec{A} \right)^2 \psi$$

$$\hat{H} = \frac{\vec{p}^2}{2m} + V(\vec{r}) \quad \rightarrow \quad \hat{H} = \frac{\left(\vec{p} - \frac{q}{c} \vec{A} \right)^2}{2m} + V(\vec{r})$$

$$\begin{aligned} \langle \vec{r} | \hat{H} | \vec{r} \rangle &= \int d^3r \psi^* \left(\frac{\vec{p}}{m} - \frac{q}{c} \vec{A} \right)^2 \psi \\ &= \int d^3r \psi^* \left(\frac{\vec{p}^2}{m} - \frac{2q}{c} \vec{p} \cdot \vec{A} + \frac{q^2}{c^2} \vec{A}^2 \right) \psi \end{aligned}$$

$$\langle \vec{r} | \hat{H} | \vec{r} \rangle = \int d^3r \frac{\langle \vec{r} | \hat{H} | \vec{r} \rangle}{\psi^* \psi}$$