

$$H = \frac{p^2}{2m} + V(r) \quad [H, J_z] = 0$$

$$\vec{L} = \vec{r} \times \vec{p}$$

$$e^{-\frac{i}{\hbar} L_z \theta} |x y z\rangle = |x' y' z'\rangle$$

$$z' = z$$

$$x' = x \cos \theta - y \sin \theta$$

$$y' = y \cos \theta + x \sin \theta$$

$$\langle \hat{n} | l m \rangle = Y_{lm}(\hat{n}) = Y_{lm}(\theta, \phi)$$

$$\psi(\vec{r}) = \underbrace{R(r)}_r \underbrace{N(\theta, \phi)}_{Y_{lm}(\theta, \phi)}$$

$$N(\theta, \phi) = Y_{lm}(\theta, \phi)$$

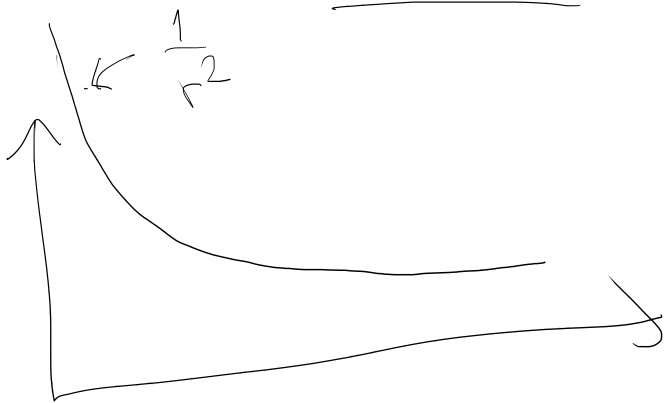
$$\frac{-\hbar^2}{2m} \frac{d^2}{dr^2} R(r) + \left[\frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} + V(r) \right] R = R(r) E$$

$$0 \leq r < \infty \quad l+1$$

$$R(r) \xrightarrow{r \rightarrow 0} 0$$

$$\psi(\vec{r}) \xrightarrow{r \rightarrow 0} 0$$

$$V_{\text{eff}}(r) = \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} + V(r)$$



Example Free particle

$$\psi = \frac{R(r)}{r} Y_{lm}(\theta, \phi)$$

$$-\frac{\hbar^2}{2m} \frac{d^2 R}{dr^2} + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} R = E R$$

$$E > 0 \quad E = \frac{\hbar^2 k^2}{2m}$$

for $l=0$

$$-\frac{\hbar^2}{2m} R'' = \frac{\hbar^2}{2m} k^2 R$$

$$\Rightarrow R = \sin(kr + \delta)$$

$$\psi = \frac{R(r)}{r} \Rightarrow R(r=0) = 0 \Rightarrow R = \sin(kr)$$

$$\psi(\vec{r}) = N \frac{\sin(kr)}{r}$$

$$E = \frac{\hbar^2 k^2}{2m}$$

For $l \neq 0$

$$-\frac{\hbar^2}{2m} \frac{d^2 R}{dr^2} + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} R = E R$$

$$R_1(r) = r J_l(kr) \quad \text{or} \quad R_2 = r n_l(kr)$$

$$R_1(r) \xrightarrow{r \rightarrow 0} r^{l+1}$$

$$R_2 \xrightarrow{r \rightarrow 0} \frac{1}{r^l}$$

$$\psi(\vec{r}) = N_{lm} j_l(kr) Y_{lm}(\theta, \phi)$$

alternative basis

$$\psi_{\vec{k}}(\vec{r}) = e^{i\vec{k}\cdot\vec{r}} = \sum_{lm} N_{lm} j_l(r) Y_{lm}(\theta, \phi)$$

$$N_{lm} = \int d^3r e^{-i\vec{k}\cdot\vec{r}} j_l(r) Y_{lm}(\theta, \phi)$$

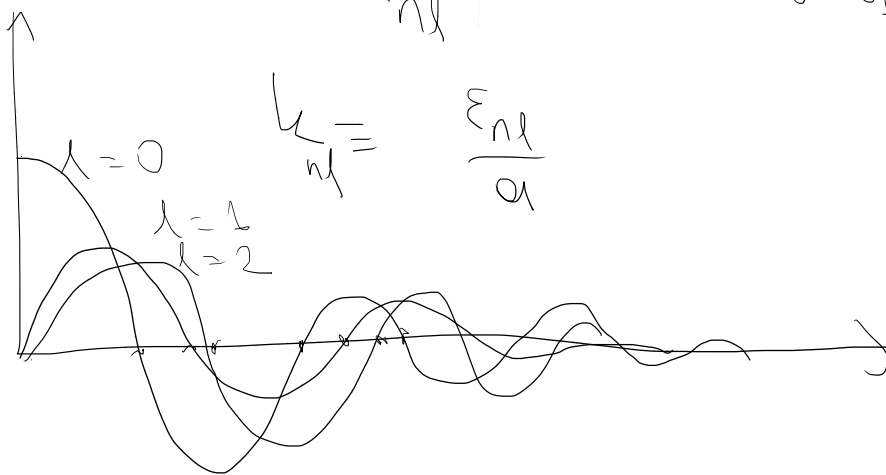
Example particle inside a ^{hard} spherical well

$$\psi(\vec{r}) = N_{lm} j_l(kr) Y_{lm}(\theta, \phi)$$

$$\psi(\vec{r})|_{r=a} = 0$$

$$\Rightarrow j_l(ka) = 0$$

$$ka = \xi_{nl} = n^{\text{th}} \text{ zero of } j_l(x)$$



Example Spherical oscillator

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 r^2$$

$$r^2 = x^2 + y^2 + z^2$$

$$H = H_x + H_y + H_z$$

$$H_i = \frac{1}{2m} p_i^2 + \frac{1}{2} m \omega^2 x_i^2$$

$$E_{n_x n_y n_z} = \hbar \omega \left(n_x + n_y + n_z + \frac{3}{2} \right)$$

$$E = \hbar \omega \left(1 + \frac{3}{2} \right)$$

$$1 = \underset{n_x}{\uparrow} 1 + \underset{n_y}{\uparrow} 0 + \underset{n_z}{\uparrow} 0 = 0 + 1 + 0 = 0 + 0 + 1$$

$$\psi = \frac{R(r)}{r} Y_{lm}$$

$$-\frac{\hbar^2}{2m} R'' + \left[\frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} + \frac{1}{2} m \omega^2 r^2 \right] R = E R$$

$$R(r) \xrightarrow{r \rightarrow \infty} r^l$$

$r \rightarrow \infty$

$$\frac{-\hbar^2}{2m} R'' + \frac{1}{2} m \omega^2 r^2 R = 0$$

$$R'' + \left(\frac{m\omega^2}{\hbar^2} \right) r^2 R = 0$$

$$\left[\frac{R}{r^2} \right]$$

$$\uparrow \left(\frac{m\omega^2}{\hbar^2} \right) = \frac{1}{a^4}$$

$r = au$ u is dimensionless

$$\frac{d^2 R}{dr^2} = \frac{1}{a^2} \frac{d^2 R}{du^2}$$

$$\frac{1}{a^2} \frac{d^2 R}{du^2} + \frac{1}{a^4} u^2 R = 0$$

$$\frac{d^2 R}{du^2} - u^2 R = 0 \quad (\text{as } u \rightarrow \infty)$$

$$\text{Hence } R = e^{-u^2/2} \quad (\text{as } u \rightarrow \infty)$$

$$E = \frac{1}{2} \hbar \omega \lambda = \frac{1}{2} \hbar \omega (2N+3)$$

$$E = \hbar \omega \left(N + \frac{3}{2} \right) \quad N \geq 1$$

For $\underline{N=1}$, $l=1 \Rightarrow \left. \begin{array}{l} m=+1 \\ m=0 \\ m=-1 \end{array} \right\}$

$$N = 0 + 1$$

$$\neq \textcircled{1} + 0 \Leftarrow$$

not even

$$N = n_x + n_y + n_z$$

$$|n_x n_y n_z\rangle$$

$$l = l(n_x, n_y, n_z)$$

$$|N l m\rangle$$

$$m = m(n_x, n_y, n_z)$$

Example Coulomb Potential

$$V(r) = -\frac{\alpha}{r} \quad \alpha = Ze^2$$

$$\psi(\vec{r}) = \frac{R(r)}{r} Y_{lm} \Leftarrow$$

$$\frac{R(r)}{r} = r^l e^{-r/a_0} w(r)$$

Example Angular momentum + spin- $\frac{1}{2}$

$$\psi(\vec{r}) = \underbrace{\psi_+(\vec{r})}_{\begin{pmatrix} 1 \\ 0 \end{pmatrix}} + \underbrace{\psi_-(\vec{r})}_{\begin{pmatrix} 0 \\ 1 \end{pmatrix}}$$
$$= \begin{pmatrix} \psi_+(\vec{r}) \\ \psi_-(\vec{r}) \end{pmatrix}$$

$$\vec{J} = \vec{L} + \vec{S} \quad \leftarrow$$

$$[L_i, S_j] = 0$$

$$Y_{lm}(\theta, \phi) \otimes \left| \pm \right\rangle$$

$$|l s; m m_s\rangle$$

$$L^2 |l s; m m_s\rangle = \hbar^2 l(l+1) |l s; m m_s\rangle$$

$$S^2 |l s; m m_s\rangle = \hbar^2 s(s+1) |l s; m m_s\rangle$$

$$L_z |l s; m m_s\rangle = \hbar m |l s; m m_s\rangle$$

$$S_z |l s; m m_s\rangle = \hbar m_s |l s; m m_s\rangle$$

$$[J_i, L_j] = [L_i + S_i, L_j] = [L_i, L_j] = i\hbar \epsilon_{ijk} L_k$$

$$[S_i, S_j] = i\hbar \epsilon_{ijk} S_k$$

$$[J_i, L^2] = L_i [S_i, L^2] + [S_i, L^2] L_i$$

$$\begin{aligned}
 &= L_j i\hbar \epsilon_{ijk} L_k + i\hbar \epsilon_{ijk} L_k L_j \\
 &= i\hbar \epsilon_{ijk} (L_j L_k + L_k L_j) = 0
 \end{aligned}$$

$$[S_i, L^2] = [L_i + S_i, L^2] = 0$$

$$[S_i, S^2] = [L_i + S_i, S^2] = 0$$

$\{J^2, J_z, L^2, S^2\}$ maximal set of commuting operators.

$$\{|l s; j m_j\rangle\}$$

$$\{|l s; m_l m_s\rangle\}$$

$$|l s; j m_j\rangle = \sum_{m_l m_s} C_{m_l m_s}^{l s j} |l s; m_l m_s\rangle$$

↑
Clebsch Gordon
Coefficients

$$C_{m_l m_s}^{l s j} = \langle l s; m_l m_s | l s; j m_j \rangle$$

$$\langle l s; m m_s | l s; j m_j \rangle = 0 \quad \text{unless } m + m_s = m_j$$

$$\langle l s; m m_s | J_z | l s; j m_j \rangle = \hbar m_j \langle l s; m m_s | l s; j m_j \rangle$$

$$\langle l s; m m_s | L_z + S_z | l s; j m_j \rangle = \hbar (m + m_s) \langle l s; m m_s | l s; j m_j \rangle$$

either $m_j = m + m_s$ or $\langle l s; m m_s | l s; j m_j \rangle = 0$

$$|l s; m m_s\rangle = |l m\rangle \otimes |s m_s\rangle$$

$$\vec{J} = \vec{L} + \vec{S} = \vec{L} \otimes \mathbb{1} + \mathbb{1} \otimes \vec{S}$$

$$J_z = L_z \otimes \mathbb{1} + \mathbb{1} \otimes S_z$$

$$L_z \otimes \mathbb{1} |l m\rangle \otimes |s m_s\rangle = (L_z |l m\rangle) \otimes (\mathbb{1} |s m_s\rangle)$$

$$L_z \otimes \mathbb{1} |l m; j m_j\rangle = \dots$$

$j \pm$ maximum possible value of m_j
 $J_{\max} =$ maximum " " " " $(m + m_s)$

$$J_{\max} = l + s$$

$$|11; \frac{3}{2}, \frac{1}{2}\rangle \propto J_- |11; \frac{3}{2}, \frac{1}{2}\rangle = \dots$$

$$|11; \frac{3}{2}, \frac{3}{2}\rangle \propto J_+ |11; \frac{3}{2}, \frac{1}{2}\rangle = \dots$$

$$\Downarrow J$$

$$|11; \frac{1}{2}, \frac{1}{2}\rangle$$

$$m_s + m = \frac{1}{2} \Rightarrow (m_s, m) = (\frac{1}{2}, 0)$$

$$\text{or}$$

$$(m_s, m) = (-\frac{1}{2}, 1)$$

$$|11; \frac{1}{2}, \frac{1}{2}\rangle = a |11; 0, \frac{1}{2}\rangle + b |11; 1, -\frac{1}{2}\rangle$$

$$\langle 11; \frac{3}{2}, \frac{1}{2} | 11; \frac{1}{2}, \frac{1}{2}\rangle = 0$$

$$= \sqrt{2}a + b = 0 \Rightarrow b = -\sqrt{2}a$$

$$|11; \frac{1}{2}, \frac{1}{2}\rangle = \frac{1}{\sqrt{3}} \left(|11; 0, \frac{1}{2}\rangle - \sqrt{2} |11; 1, -\frac{1}{2}\rangle \right)$$

$$|11; \frac{1}{2}, \frac{1}{2}\rangle \propto J_- |11; \frac{1}{2}, \frac{3}{2}\rangle$$

$$|J_1 J_2; m_1 m_2\rangle$$

$$|J_1 - J_2| \leq J \leq J_1 + J_2$$

$$\begin{aligned} & \left\{ \begin{aligned} & |J_1 J_2; \underbrace{J_1 + J_2}_{m_1} \underbrace{J_1 + J_2}_{m_2}\rangle = |J_1 J_2; J_1 + J_2\rangle \\ & |J_1 J_2; J_1 + J_2, J_1 + J_2 - 1\rangle \dots |J_1 J_2; J_1 + J_2, J_1 + J_2 - J\rangle \\ & \vdots \\ & |J_1 J_2; J_1 + J_2, -J_1 - J_2\rangle \end{aligned} \right. \end{aligned}$$

$$\left\{ |J_1 J_2; J_1 + J_2 - 1, J_1 + J_2 - 1\rangle \right.$$

$$= a |J_1 J_2; J_1 + J_2 - 1\rangle + b |J_1 J_2; J_1 - 1, J_2\rangle$$

$$\boxed{|J_1 - J_2| \leq J \leq J_1 + J_2}$$